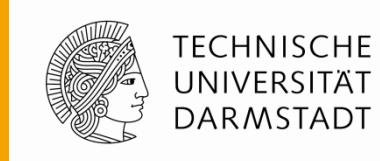


Self-consistent model of charge-carrier injection: extension to the case of a degenerate organic semiconductor

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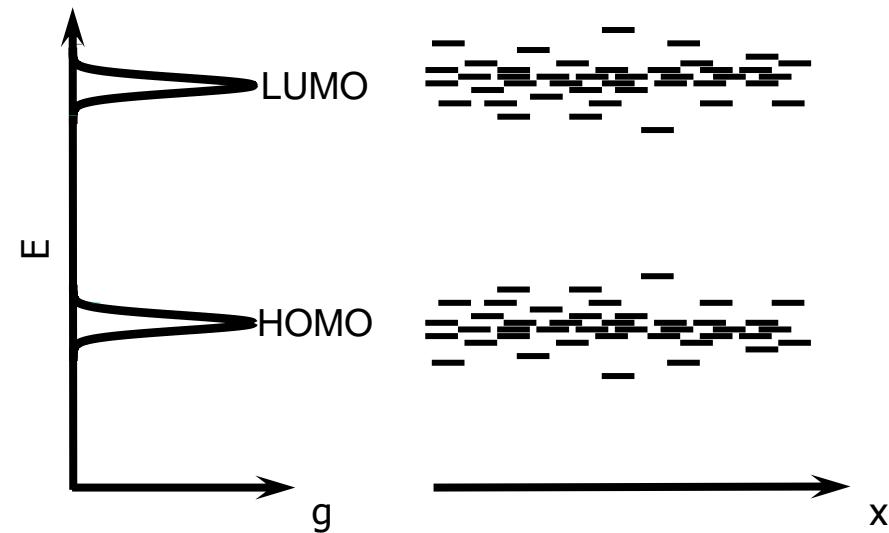
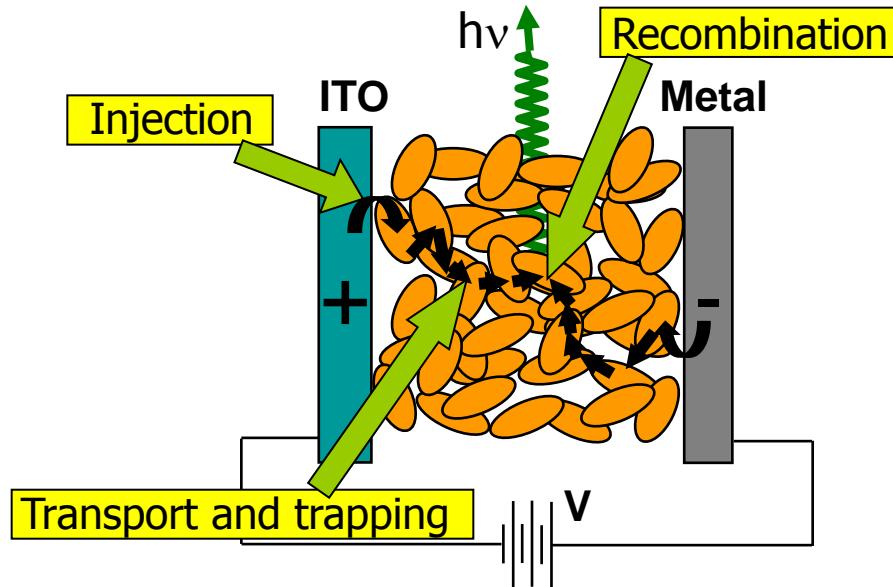
Outline



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- Outline of the previous self-consistent model of charge-carrier injection and transport accounting for discrete nature of charge carriers (MMF approach)
- Approximations for the charge-carrier density in the case of Gaussian DOS
- Extension of the MMF model to the case of degenerate semiconductor: voltage-barrier chart, current-voltage characteristics for unipolar devices
- Conclusion

Important factors of electronic processes in organic semiconductors and our aims



- We need in a **self-consistent model of charge injection and transport** in organic semiconductors allowing to extract from experiments important **device parameters** (carrier concentrations, mobilities, injection barriers, field distribution) and to **evaluate the fatigue stability** of specific devices

Self-consistent 1D mean-field model of charge-carrier injection and unipolar transport in a single-layer device



Keypoints of the model

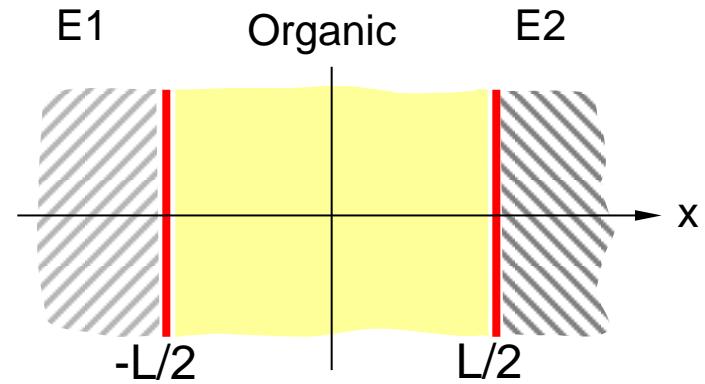
A. Organic semiconductor:

$$kT \mu_i \frac{dp_i(x)}{dx} - e\mu_i p_i(x) F_i(x) = -j$$

$$\frac{dF_i(x)}{dx} = \frac{e}{\epsilon_i \epsilon_0} p_c(x)$$

narrow DOS: $g(E) = P_c \delta(E - E_c)$

Boltzmann statistics of the carriers: $p_i(x) = P_c \exp \left[\frac{E_c - \kappa_i(x) - e\varphi(x)}{kT} \right]$



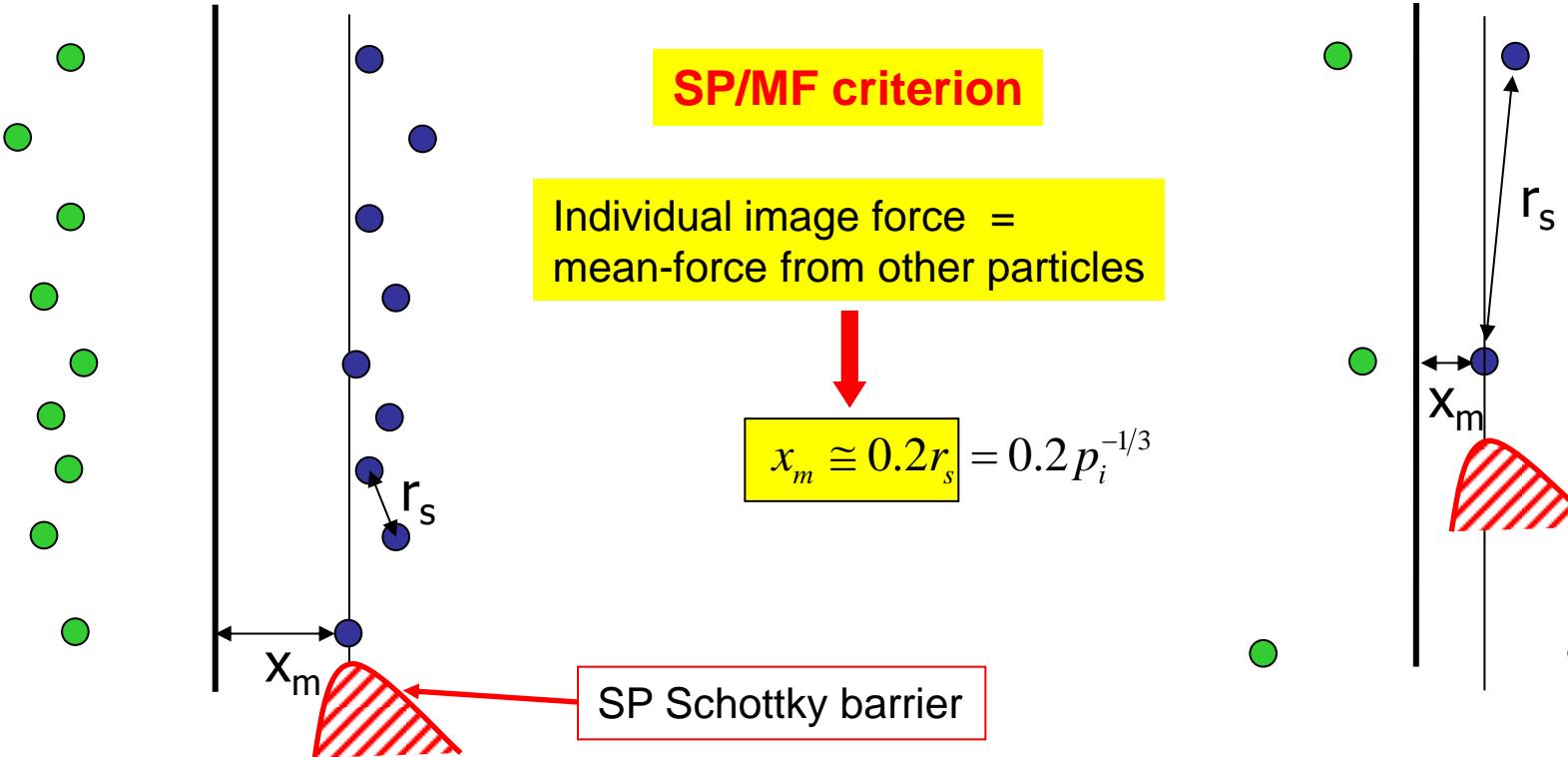
B. Electrodes: Thomas-Fermi approximation

C. Self-consistent boundary conditions at the electrode/OSC interfaces:

- electric displacement and electrochemical potential are continuous
- electrostatic potential is continuous (mean-field (MF) approach) or has a jump $\Delta\varphi_e = \delta\varphi_{Sch}$ (modified mean-field (MMF) model)

Neumann et al. JAP **100**, 084511 (2006); Genenko et al., PRB **81**, 125310(2010)

Charge-carrier injection: many-particle vs single-particle (SP) mechanism

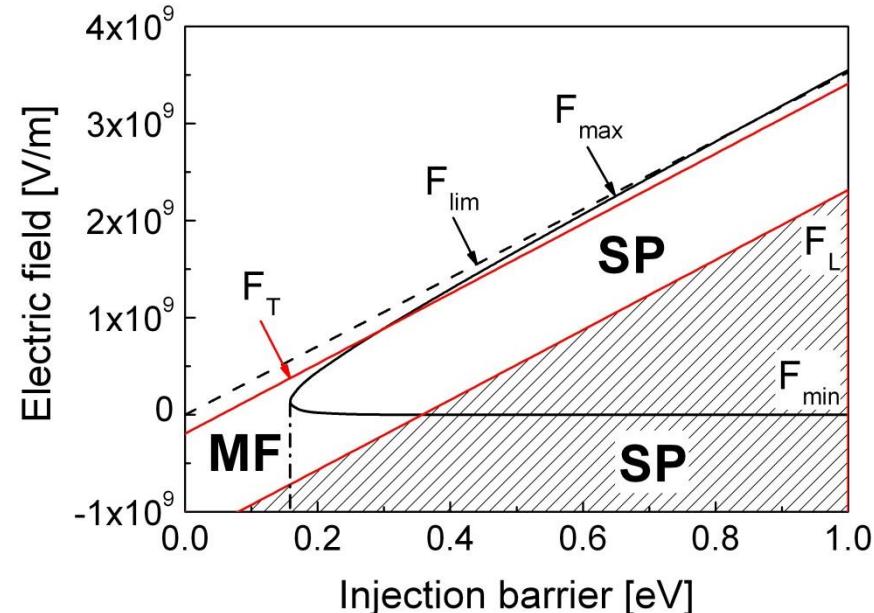
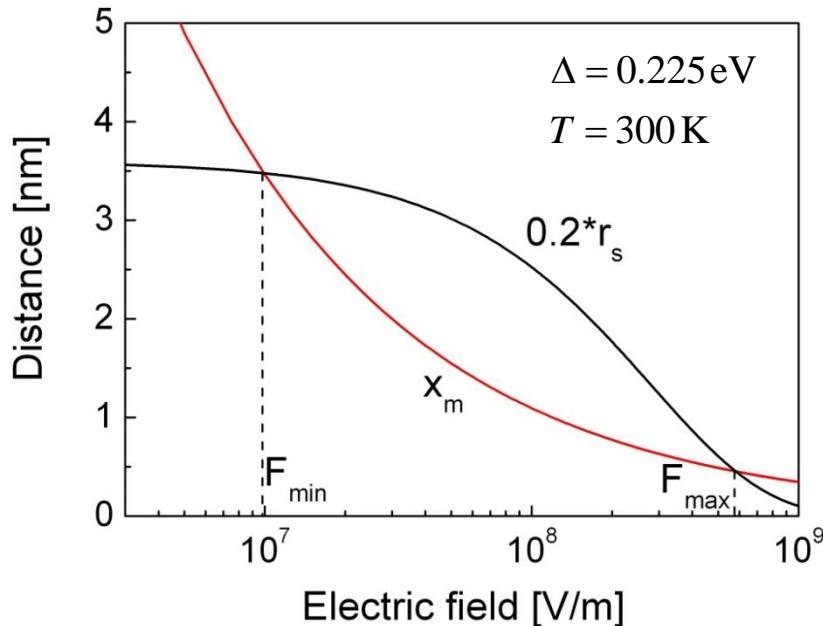


$$x_m^\pm = \sqrt{\frac{e}{16\pi\varepsilon_0\varepsilon_i |F_i(\pm L/2)|}}$$

$$p_i(\pm L/2) = P_c \exp \left\{ -\frac{\Delta^\pm}{kT} \mp \frac{el_{TF}^\pm}{kT} \left[\frac{\varepsilon_i}{\varepsilon_e^\pm} F_i(\pm L/2) - \frac{i}{V_e^\pm} \right] \right\}$$

Genenko et al., PRB **81**, 125310(2010)

Validity of SP and MF approximations: field region



ITO	Organic
$n_\infty^- = 10^{20} \text{ cm}^{-3}$	$P = 10^{21} \text{ cm}^{-3}$
$\epsilon_e^- = 9.3$	$\epsilon_i = 3$
$\mu_e^- = 30 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$	$\mu_p = 10^{-4} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$
$\kappa_\infty^- = 0.225 \text{ eV}$	

Additionally:

$$r_s = L \Rightarrow F_L$$

$$r_s = l_{Coulomb} = e^2 / 32\pi\epsilon_0\epsilon_i kT \Rightarrow F_T$$

Account of charge-carrier discreteness in the MF approach

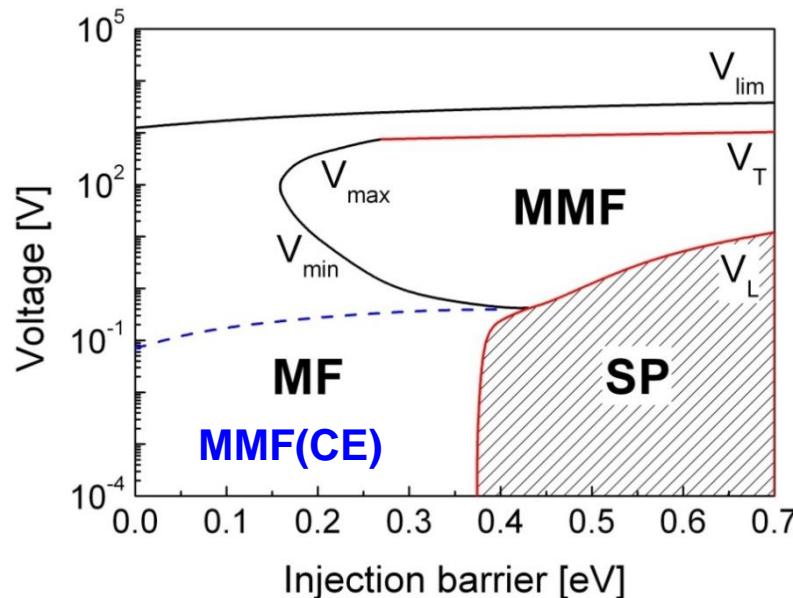


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Boundary conditions:

$$p_i(\pm L/2) = P_c \exp \left[-\frac{\Delta^\pm}{kT} m \frac{eF_i(\pm L/2)l_{TF}^\pm}{kT} \frac{\varepsilon_i}{\varepsilon_e^\pm} + \left(1 - \frac{x_m^\pm}{0.2r_s^\pm} \right) \frac{e\delta\varphi_{sch}^\pm}{kT} \theta(0.2r_s^\pm - x_m^\pm) \right]$$

ITO/Organic/AI



$$V = \int_{-\infty}^{\infty} dx F(x) - V_{BI}$$

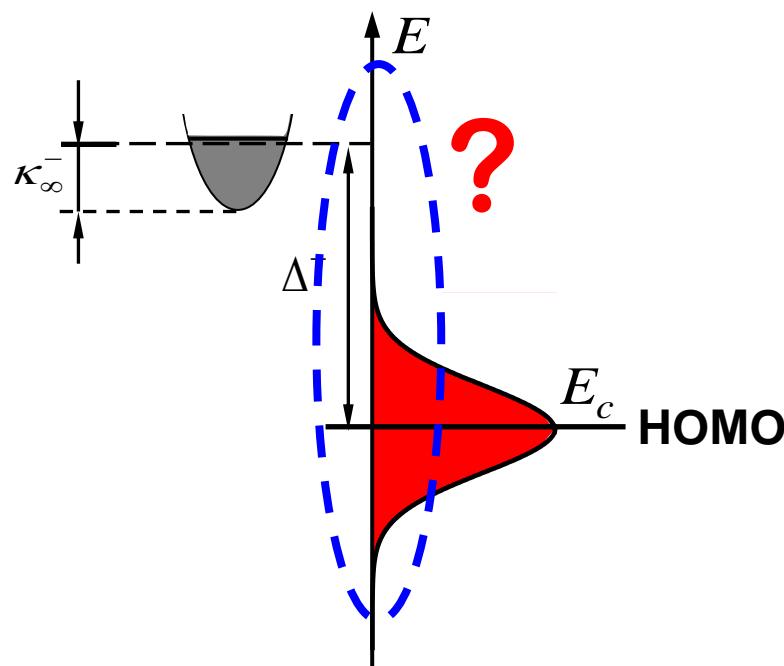
$$V_{BI} = \int_{-\infty}^{\infty} dx F(x) \Big|_{j=0}$$

Genenko et al., PRB **81**, 125310(2010)

What changes for the Gaussian DOS?



ITO/OSC interface



$$\begin{aligned} p_i(x) &= \int_{-\infty}^{\infty} dE g(E) f(E, \kappa_i(x)) \\ &= \frac{P_c}{\sqrt{2\pi}\sigma_c} \int_{-\infty}^{\infty} dE \frac{\exp\left(-\frac{(E-E_c)^2}{2\sigma_c^2}\right)}{\exp\left(\frac{\kappa_i(x)+e\varphi(x)-E}{kT}\right)+1} \end{aligned}$$

In the „Boltzmann regime“:

$$\Delta^\pm \rightarrow \Delta^\pm - \frac{\sigma_c^2}{2kT}$$

- Whether the Boltzmann statistics of carriers is still relevant in the case of realistic Gaussian DOS of OSC, especially for calculation of the carrier density at the electrode/OSC interfaces?

Charge-carrier density as Gauss-Fermi integral

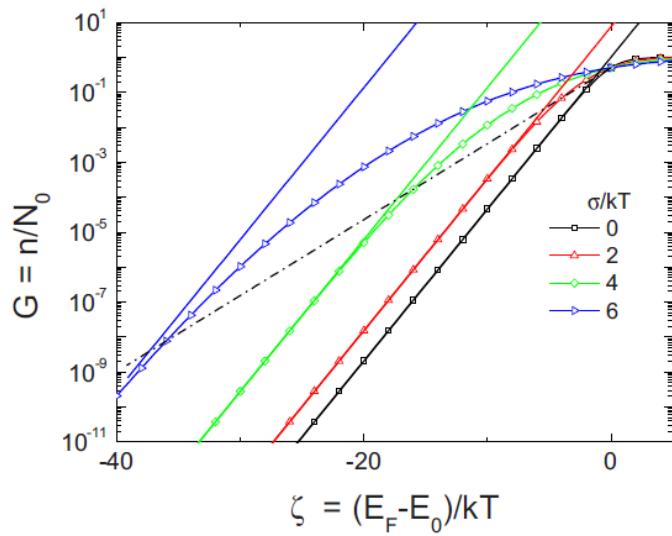


Paasch & Scheinert, JAP **107**, 104501 (2010)

$$n = \int dE g(E) f(E, E_F) = \frac{N_0}{\sqrt{2\pi}\sigma} \int_{-\infty}^{\infty} dE' \exp\left(-\frac{E'^2}{2\sigma^2}\right) \left[\exp\left(\frac{E' - E_F + E_0}{kT}\right) + 1 \right]^{-1}$$

Nondegenerate limit: $f(E, E_F) \rightarrow \exp\left(\frac{E_F - E}{kT}\right)$

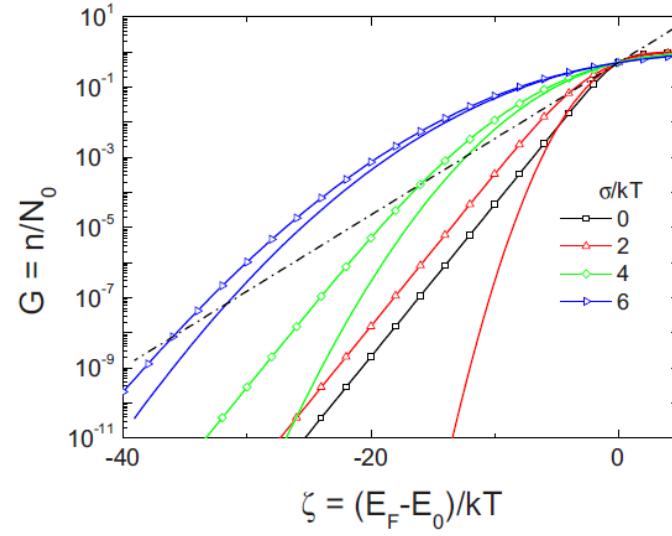
$$n = N_0 \exp\left(\frac{E_F - E_0}{kT} + \frac{\sigma^2}{2(kT)^2}\right)$$



Crossover at $\zeta = -\frac{\sigma^2}{2(kT)^2}$: $n_{cr} = \frac{N_0}{2} \exp\left(-\frac{\sigma^2}{2(kT)^2}\right)$

Degenerate limit: $f(E, E_F) \rightarrow \theta(E_F - E)$

$$n = \frac{N_0}{2} \operatorname{erfc}\left(-\frac{E_F - E_0}{\sigma\sqrt{2}}\right)$$



For $\sigma(300K) = 0.1$ eV $n_{cr} \sim 10^{-4} N_0$



PS analytical approximation for GF integral

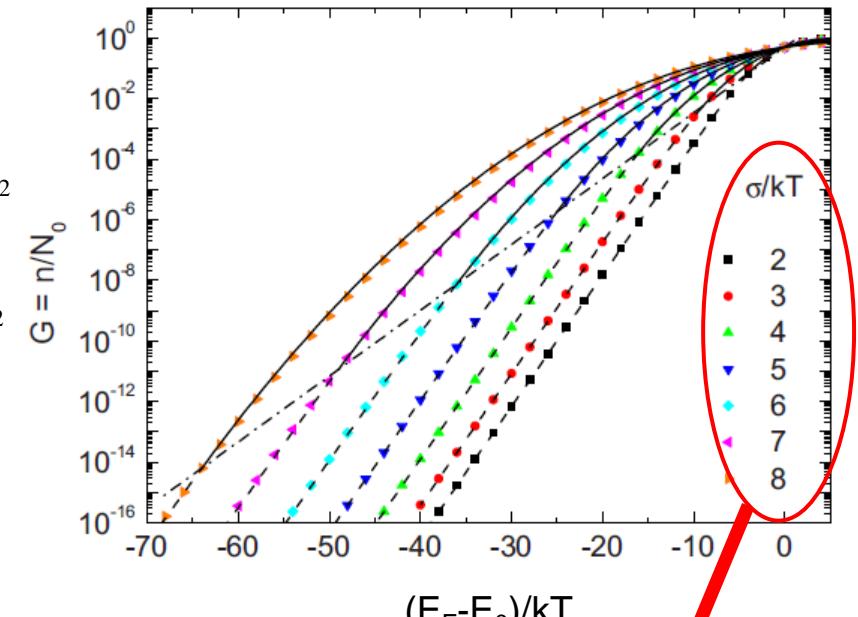
Paasch & Scheinert, JAP 107, 104501 (2010)

$$\zeta = \frac{E_F - E_0}{kT}, \quad s = \frac{\sigma}{kT}$$

$$\frac{n}{N_0} = \begin{cases} \exp\left(\zeta + \frac{s^2}{2}\right) \left[\exp\left(K(s)(\zeta + s^2)\right) + 1 \right]^{-1}, & \zeta \leq -s^2 \\ \frac{1}{2} \operatorname{erfc}\left(-\frac{\zeta}{s\sqrt{2}} H(s)\right), & \zeta \geq -s^2 \end{cases}$$

$$H(s) = \frac{\sqrt{2}}{s} \operatorname{erfc}^{-1} \left[\exp\left(-\frac{s^2}{2}\right) \right]$$

$$K(s) = 2 \left\{ 1 - \frac{H}{s} \sqrt{\frac{2}{\pi}} \exp\left[\frac{s^2}{2}(1-H^2)\right] \right\}$$



$$\sigma_c(300K) = 0.05 - 0.2 \text{ eV}$$

- After appropriate modifications this approximation can be applied to the MMF boundary conditions

MMF boundary conditions in the case of Gaussian DOS



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Self-consistent boundary conditions at the electrode/OSC interfaces are:

$$p_i\left(\pm \frac{L}{2}\right) = \frac{P_c}{\sqrt{2\pi}\sigma_c} \int_{-\infty}^{\infty} dE \exp\left(-\frac{E^2}{2\sigma_c^2}\right) \left[\exp\left(\frac{\Delta_{eff}^{\pm} - E}{kT}\right) + 1 \right]^{-1}$$

$$\Delta_{eff}^{\pm} = \Delta^{\pm} \pm eI_{TF}^{\pm} \left(\frac{\epsilon_i^{\pm}}{\epsilon_e^{\pm}} F(\pm L/2) - \frac{j}{\gamma_e^{\pm}} \right) - \left(1 - \frac{x_m^{\pm}}{0.2r_s^{\pm}} \right) e\delta\varphi_{sch}^{\pm} \theta(0.2r_s^{\pm} - x_m^{\pm})$$

With substitutions $\zeta\left(\pm \frac{L}{2}\right) = \frac{\Delta_{eff}^{\pm}}{kT}$, $s = \frac{\sigma_c}{kT}$:

$$\frac{p_i(\pm L/2)}{P_c} = \begin{cases} \exp\left(\zeta\left(\pm \frac{L}{2}\right) + \frac{s^2}{2}\right) \left[\exp\left(K(s)\left(\zeta\left(\pm \frac{L}{2}\right) + s^2\right)\right) + 1 \right]^{-1}, & \zeta\left(\pm \frac{L}{2}\right) \geq s^2 \\ \frac{1}{2} \operatorname{erfc}\left(-\frac{\zeta(\pm L/2)}{s\sqrt{2}} H(s)\right), & \zeta\left(\pm \frac{L}{2}\right) \leq s^2 \end{cases}$$

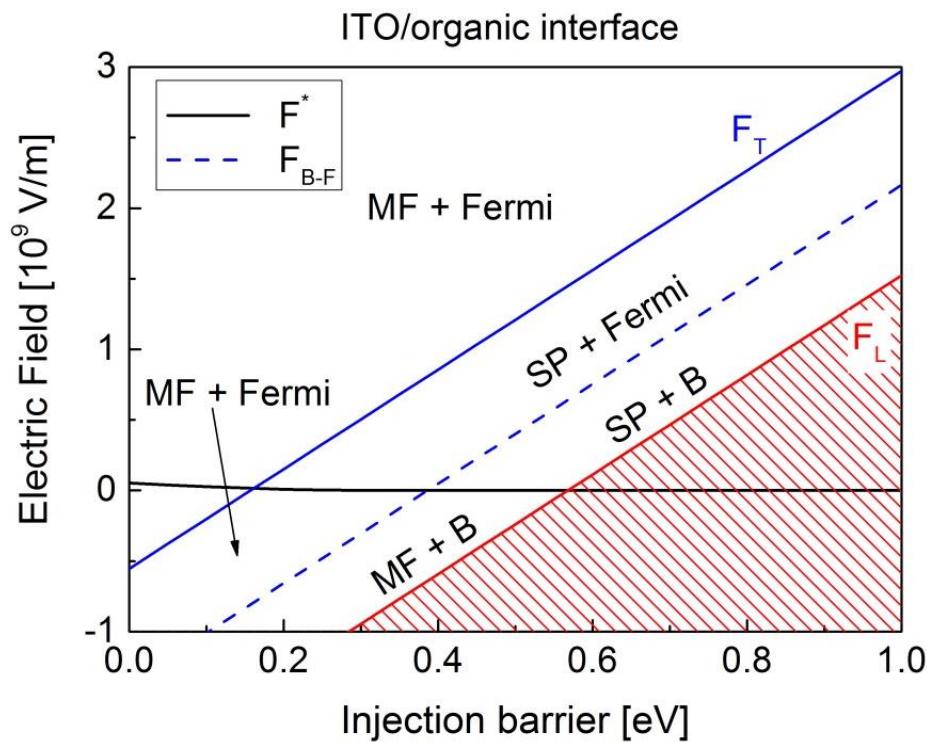
Field-injection barrier chart



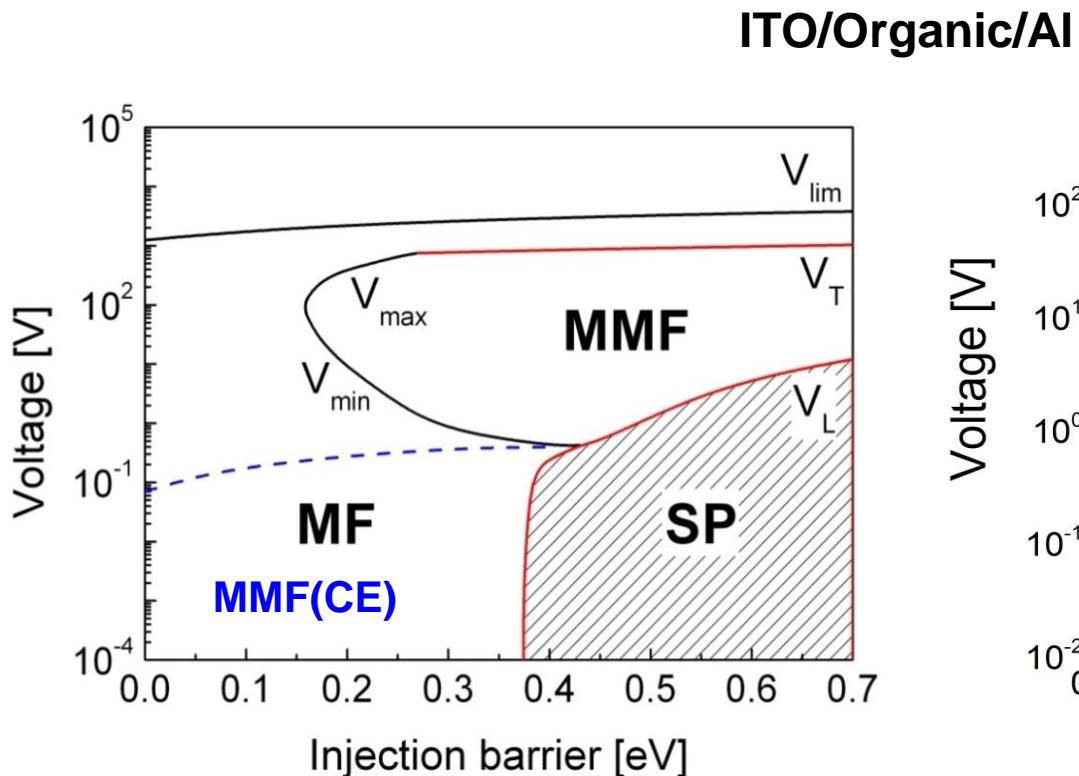
$$x_m = 0.2 p_i^{-1/3} \Rightarrow F^*$$

$$\Delta^- - e l_{TF}^- \left(\frac{\varepsilon_i^-}{\varepsilon_e^-} F(-L/2) - \frac{j}{\gamma_e^-} \right) = \frac{\sigma_c^2}{kT} \Rightarrow F_{B-F}$$

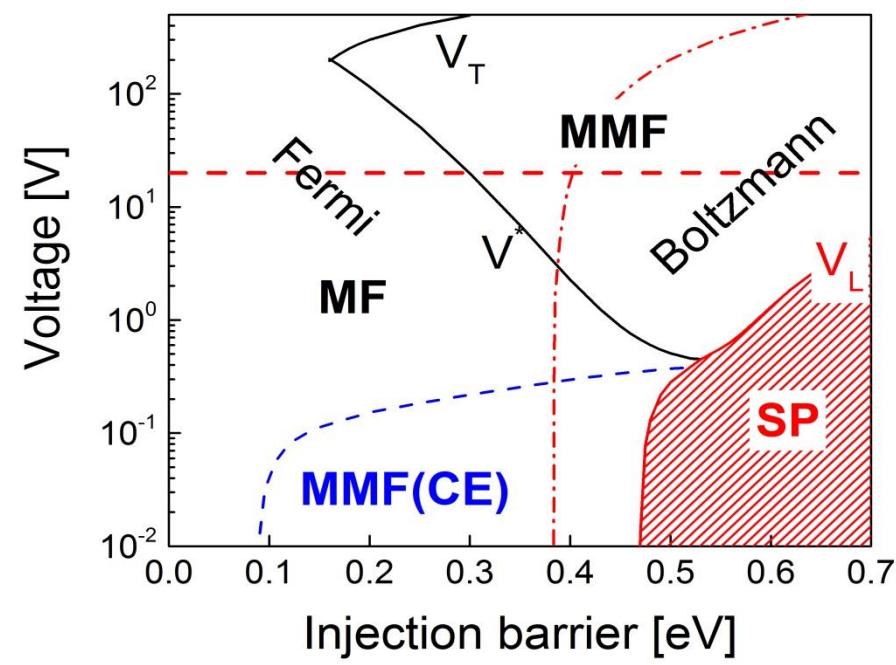
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$\mu_e^- = 30 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$	$\mu_p = 10^{-6} \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$
$\kappa_\infty^- = 0.225 \text{ eV}$	$\sigma_c = 0.1 \text{ eV}$



Voltage-barrier chart for injecting electrode



Boltzmann statistics only

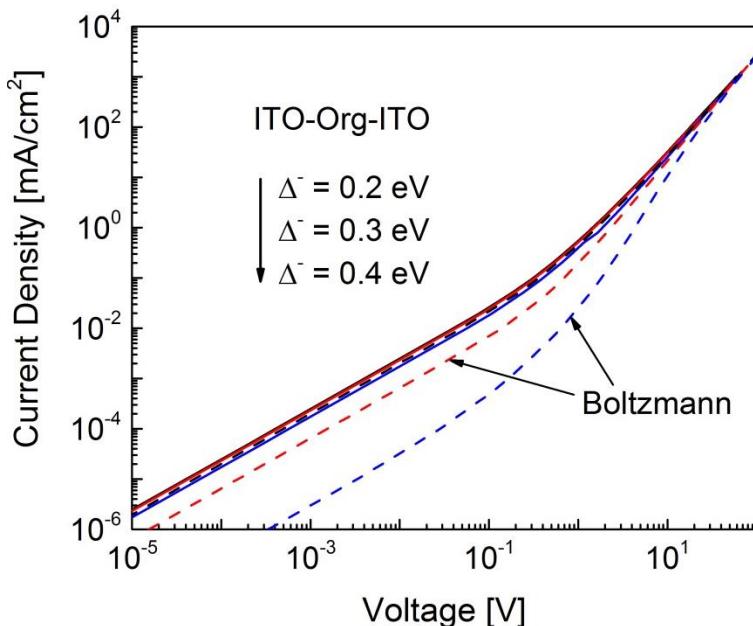


extended approach

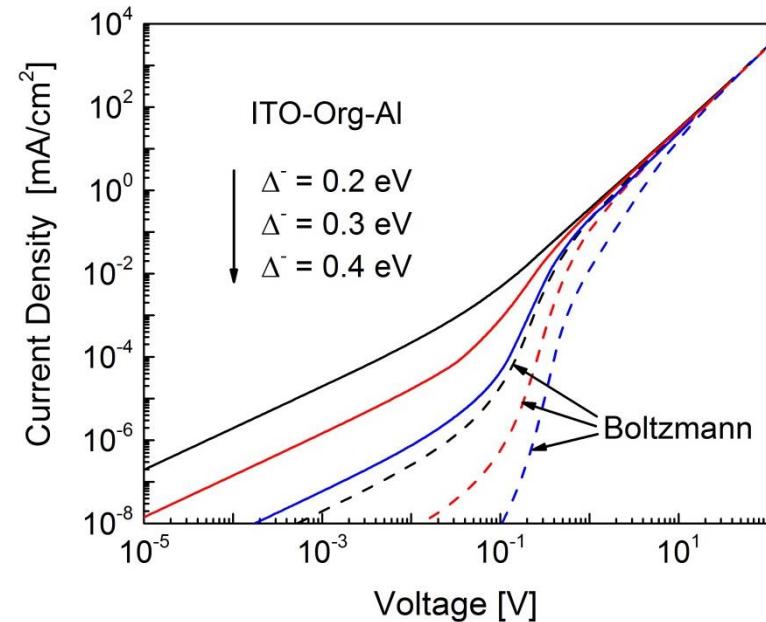
Simulated current-voltage characteristics



ITO/OSC(100nm)ITO



ITO/OSC(100nm)Al



- Assumption of the „Boltzmann regime“ for injected carriers results in sufficient undervalueing of calculated current density, especially at low voltages

Conclusions



- The mean-field, self-consistent model of charge-carrier injection and transport in organic semiconductors accounting for discrete nature of charge carriers is extended to the case of degenerate OSC with the Gaussian shape of DOS
- The extended approach is applicable now for arbitrary values of injected carrier densities in the wide range of DOS parameters
- It is expected to improve the fitting of measured current-voltage characteristics, particularly at applied voltages near and well below the built-in potential value
(Also, for fitting of fatigued bipolar IV characteristics –
the challenge unresolved yet...)